



The impact of an extreme solar event on the middle atmosphere: a case study

Thomas Reddmann¹, Miriam Sinnhuber¹, Jan Maik Wissing², Olesya Yakovchuk², and Ilya Usoskin³

¹Institute of Meteorology and Climate Research, Karlsruhe Institute of Technology, Karlsruhe, Germany

²Institute for Physics, University of Rostock, Rostock, Germany

³Space Physics and Astronomy Research Unit and Sodankyla Geophysical Observatory,
University of Oulu, Oulu, Finland

Correspondence: Thomas Reddmann (thomas.reddmann@kit.edu)

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Abstract. A possible impact of an extreme solar particle event (ESPE) on the middle atmosphere is studied for present-day climate and geomagnetic conditions. We consider an ESPE with an occurrence probability of about 1 per millennium. In addition, we assume that the ESPE is followed by an extreme geomagnetic storm (GMS), and we compare the contribution of the two extreme events. The strongest known and best-documented ESPE of 774/5 CE is taken as a reference example and established estimates of the corresponding ionization rates are applied. The ionization rates due to the energetic particle precipitation (EPP) during an extreme GMS are upscaled from analyzed distributions of electron energy spectra of observed GMSs. The consecutive buildup of NO_x and HO_x by ionization is modeled in the high-top 3D chemistry circulation model Karlsruhe Simulation Model of the middle Atmosphere (KASIMA), using specified dynamics from ERA-Interim analyses up to the stratopause. A specific dynamical situation was chosen that includes an elevated stratosphere event during January and maximizes the vertical coupling between the northern polar mesosphere–lower thermosphere region and the stratosphere; it therefore allows us to estimate a maximum possible impact. The particle event initially produces about 65 Gmol of NO_y , with 25 Gmol of excess NO_y even after 1 year. The related ozone loss reaches up to 50 % in the upper stratosphere during the first weeks after the event and slowly descends to the mid-stratosphere. After about 1 year, 20 % ozone loss is still observed in the northern stratosphere. The GMS causes strong ozone reduction in the mesosphere but plays only a minor role in the reduction in total ozone. In the Southern Hemisphere (SH), the long-lived NO_y in the polar stratosphere, which is produced almost solely by the ESPE, is transported into the Antarctic polar vortex, where it experiences strong denitrification into the troposphere. For this special case, we estimate a NO_3 washout that could produce a measurable signal in ice cores. The reduction in total ozone causes an increase of the UV erythema dose of less than 5 %, which maximizes in spring for northern latitudes of 30° and in summer for northern latitudes of about 60°.

1 Introduction

Strong events of solar activity, such as the Carrington flare event (Carrington, 1859), have been observed for more than 150 years; since the satellite era, their energetic output has been well documented (see, for example, Aschwanden et al., 2017). These events and connected strong geomagnetic storms (GMSs) often produce a strong flux of energetic particles which, when precipitating into the Earth's at-

mosphere, by ionization processes (Usoskin et al., 2011), cause the formation of the radicals NO_x and HO_x . This results in an additional ozone destruction in the middle atmosphere (see, for example, Sinnhuber et al., 2012). For solar eruptions observed in recent decades, there are many studies, observational and modelling, that have documented this solar–terrestrial interaction. As an example, see studies in the context of the World Climate Research

Programme's Stratosphere-troposphere Processes And their Role in Climate (SPARC) High Energy Particle Precipitation in the Atmosphere (HEPPA) initiative such as Funke et al. (2011a, 2017) and Sinnhuber et al. (2021). In addition, some studies indicate a possible small but not conclusive surface impact (see, for example, Seppälä et al., 2009; Maliniemi et al., 2014; Calisto et al., 2011). Even stronger events with a probability of $\sim 1/1000 \text{ yr}^{-1}$ have marked the concentration of cosmogenic isotopes in natural stratified records (e.g., tree trunks, ice cores) during the last millennia (Usoskin, 2017; Cliver et al., 2022). The impact of such extreme events has been studied mostly in a historical context in order to simulate the production of isotopes in the lower stratosphere and their transport to the surface to compare with the isotopic records (e.g., Sukhodolov et al., 2017). Here we apply ionization rates estimated for an extreme solar event in our chemical model of the atmosphere to study its chemical impact and its consequences for the ozone layer under present-day conditions, i.e., if such an event occurred in the present. The main objective of this study is to estimate the maximum direct impact which such an extraordinary event may have on the chemical state of the atmosphere and consequentially for the on-ground erythemal UV dose. As ozone-destroying radicals are generated over a wide range of altitudes in the middle atmosphere, the specific dynamical situation also determines its impact. We focus on the Northern Hemisphere (NH) and apply a dynamical situation where the vertical coupling in the middle atmosphere is particularly efficient.

The paper is structured as follows: first, we construct an extreme solar scenario with properties of an extreme event observed in the past or estimated to have an occurrence rate of about $\sim 1/1000 \text{ yr}^{-1}$ and estimate the corresponding ionization rate. Then we apply this scenario to the specific atmospheric situation and calculate the additional NO_x and its impact on ozone. Finally, we estimate the consequences for the additional UV dose following the event for NH mid-latitudes.

2 Extreme scenario setup

A solar eruption event has several components: the very first signal of such an event detectable on Earth is a flare of electromagnetic radiation emitted in the eruption at the solar surface. Observable from γ ray energies down to the visual spectrum (white flare), the total impact of the flare on the Earth's middle atmosphere is small and negligible in terms of the chemical impact (Pettit et al., 2018) and therefore in terms of possible dynamical coupling. Then, particles accelerated in this explosion and farther in the corona and interplanetary medium reach the Earth within minutes to hours after the flare, known as the solar proton event (SPE). Here the highest energies of the particles are observed up to several giga-electron volts (GeV), which initiate a nucleonic cascade in the atmosphere. This can reach even the ground and, when colliding with atmosphere's main constituents,

cause nuclear reactions which can form the cosmogenic isotopes ^{14}C or ^{10}Be . As one of the strongest of such events detected in paleo-nuclide records so far, the possible solar eruption dated at 774/5 CE has been studied extensively (see, e.g., Usoskin, 2017), including its impact on the atmosphere (Sukhodolov et al., 2017). Compared to the strongest directly observed SPE in 1956, Mekhaldi et al. (2021) estimated a factor 40 higher fluence as having produced the observed nuclide concentrations. More recent studies hint at an even stronger fluence. Cliver et al. (2022) estimate a $70\times$ particle fluence compared to the 1956 event, and we scale the ionization rates of the 1956 event accordingly, viz. $70\times$.

In addition, a strong eruption on the Sun's surface is often accompanied by a "coronal mass ejection", which, when directed to Earth, hits the magnetosphere and causes major GMSs, typically in the range of 2 d after the flare event. The energies of the particles of the GMS (mostly electrons) are in the range of a few kilo-electron volts (keVs) to about 1 MeV, and they impact mainly the mesosphere and lower thermosphere. The strength of the disturbance of the geomagnetic field can be expressed by planetary indices like Kp and Ap, but their relation to ionization rates for extreme cases can be highly non-linear. Here we take a different approach and assume the results of the study of Meredith et al. (2016) to be applicable for strong geomagnetic events when studying the interaction with the atmosphere. Meredith et al. (2016) conducted an extreme value analysis of the electron flux measured in the MEPED (Medium Energy Proton and Electron Detector) 90° pitch angle detector for the three energy intervals provided by the instrument. We assume here that we can apply these results also for particles entering the atmosphere (0° pitch-angle) and scale the observed strong GMS events according their distribution. We select extreme events at 30 and 300 keV for the values 4.5 and 6 of the parameter L_* (Meredith et al., 2016, Tables 3 and 4) and take the strongest events as representatives of extreme geomagnetic events. L_* is related to the geomagnetic latitude where the particles enter the atmosphere, which is normally the auroral oval ($L_* = 6$), but for disturbed conditions it can be slightly shifted to lower latitudes ($L_* = 4.5$). Note that their list excludes SPEs. The 20 November 2003 event is the highest at $L_* = 4.5$ and third in rank at $L_* = 6$ for 30 keV, the 6 April 2010 09:00–12:00 UT event is highest at $L_* = 6$, and the several 6 April 2010 events are together highest also at $L_* = 4.5$ for 300 keV. The electron fluxes for these events have a probability of $\sim 1/10 \text{ yr}^{-1}$ when inspecting their Fig. 6. From this figure, one can also deduce enhancement factors for $1/100 \text{ yr}^{-1}$ events which are about $2\times$ for low- and mid-energy electrons and $10\times$ for high-energy electrons compared to the $1/10 \text{ yr}^{-1}$ events. Our extreme GMS event is finally constructed in the following way: we take the ionization rates (including protons) calculated with the AIMOS-AISstorm 2.0 model (described in Nesse Tyssøy et al., 2021), interpolate them to our model grid for 20 November 2003 and 6 April 2010, scale them with factors 2 and 10, respec-

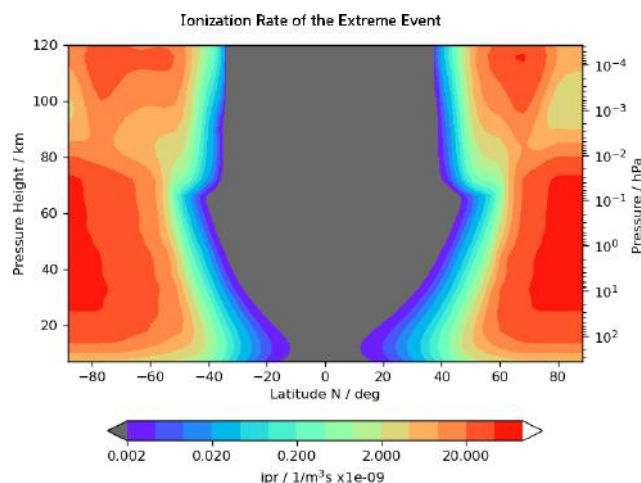


Figure 1. The ionization rate of the extreme event (mean of 23–25 January 2009). One can clearly see the GMS component maximizing in the lower thermosphere following the auroral oval and the polar cap ionization by the SPE.

tively, and add them up as one event with a duration of 1 d. As the AIMOS-AISstorm model includes electrons limited to < 300 keV, the lower boundary of a realistic electron ionization rate is about 70 km. The role of an extreme GMS in the NO_y budget has not been studied before. To assess its impact on the ozone layer, high-top models are necessary for calculation of NO_y in the source region and its subsequent downward transport.

Finally, we assume that SPE and GMS indeed can be combined for this sensitivity study, starting with a 1 d extreme SPE and after 2 d with the extreme GMS. Figure 1 shows the 3 d and zonal average ionization rate. We note that this approach may not be completely physically consistent but should yield a meaningful upper value for the strength of an extreme event.

The impact of the solar events on ozone is mainly given by the particle-induced buildup of NO_x which catalytically reacts with ozone. NO_y , which includes all nitrogen containing species except N_2O including the reservoir gases for NO_x , has a very long lifetime in the stratosphere, where it contributes effectively to ozone destruction. In the middle atmosphere, the impact on ozone is the key parameter for any dynamical atmospheric impact of such events, as short-wave heating rates are dominated by ozone. Changes in ozone will change the temperature distribution in the middle atmosphere and the winds and propagating waves as a result. In addition, as the ozone density maximizes in the stratosphere, any reduction of ozone will also change the UV intensity at the surface. The SPE-induced ionization rate maximizes in the upper stratosphere and lower mesosphere; the GMS has its maximum ionization in the lower thermosphere. In order to estimate a possible maximum impact of a particular solar event, an appropriate dynamical situation would be where the

portion of NO_y built up in the mesosphere–lower thermosphere (MLT) during the event survives photochemical destruction and is transported into the lower stratosphere, where its lifetime is on the order of many years. This sets the date of such an experiment to the winter season.

In satellite observations from instruments like the Michelson Interferometer for Passive Atmospheric Sounding (MIPAS) on Envisat (Fischer et al., 2008) strong NO_y intrusions from the lower thermosphere into the NH mesosphere and upper stratosphere have been observed after mid-winter sudden stratospheric warmings (SSWs) when accompanied by elevated stratosphere events. These show a strong downward transport of air from the MLT (Holt et al., 2013). For the NH, it has been shown that these events also yield the strongest ozone impact (Sinnhuber et al., 2018). In order to test if an extreme GMS can have a similar impact as the SPE, which maximizes lower in the mesosphere, we synchronize the SSW and the GMS. As an example of such a dynamical situation, we use the SSW of 21 January 2009, which has also been studied in the HEPPA-II experiment (Funke et al., 2017). We note that the situation in the Southern Hemisphere (SH) may be different. Here the polar middle atmosphere is essentially undisturbed for the whole winter, and an early intrusion from the MLT region may reach the mid-stratosphere. This has also been observed for the Antarctic winter 2003 (Funke et al., 2005).

In summary, the extreme event we study is a combination of a SPE and a GMS. The ionization rate is derived from observed events which have been scaled to a strength estimated for an event with a occurrence rate of $\sim 1/100 \text{ yr}^{-1}$ to $\sim 1/1000 \text{ yr}^{-1}$, synchronized with an elevated stratosphere event.

3 Description of the model and experiment

For the purpose of this sensitivity study, we use the Karlsruhe Simulation Model of the middle Atmosphere (KASIMA; Kouker et al., 1999) in the version described in Sinnhuber et al. (2021). The model solves the meteorological basic equations in spectral form in the altitude range between 300 hPa and 3.6×10^{-5} hPa, with pressure height $z = H \log(p/p_0)$ ($H = 7$ km and $p_0 = 1013.25$ hPa) as a vertical coordinate. It uses radiative forcing terms for UV–Vis and IR as well as a gravity wave drag scheme. In order to yield a realistic meteorology, the model is relaxed (nudged) to ERA-Interim meteorological analyses (Dee et al., 2011) between the lower boundary of the model and 1 hPa. A full stratospheric chemistry including heterogeneous processes is adapted to include source terms related to particle ionization. For the production of HO_x the parameterization of Solomon et al. (1981) is used. For the production of NO_x , 0.7 NO molecules are produced per ion pair and 0.55 N atoms in ground state. The model also includes a HNO_3 production from proton hydrates based on the parameterization of

de Zafrá and Smyshlyaev (2001), which has been modified to be dependent on actual ionization rates. The model was used in all the three HEPPA inter-comparisons (Funke et al., 2011a, 2017; Sinnhuber et al., 2021). The model has proven to realistically simulate chemistry in the middle atmosphere for NO_y intrusions, which is sufficient for studying its direct impact.

The following simulations are performed: one with the extreme event (EXT), where the ionization rates from the extreme event are applied from 21 to 23 January, and one with a background ionization only (REF). The background ionization rate is a mean from AIMOS ionization rates (Wissing and Kallenrode, 2009) for minimum Ap indices. In addition, a simulation which only includes the solar proton event is used to compare the contributions of the two components.

4 Results

4.1 NO_y intrusion

Figure 2 shows the mean of NO_y ($=\text{NO}_x + \text{HNO}_3 + \text{ClONO}_2 + \text{BrONO}_2 + \text{HONO} + \text{N}_2\text{O}_5 + \text{BrONO}$; gas phase only), where $\text{NO}_x = \text{N} + \text{NO} + \text{NO}_2$, over the northern polar cap (latitude $> 70^\circ \text{N}$) for 1 year with the dynamics and chemical boundary conditions of 2009 but with the extreme solar event described in Sect. 2 included. The two components of the event (SPE and GMS) can clearly be seen as an instantaneous increase over all altitudes and a tongue of increased NO_y values transported from the MLT within the elevated stratopause event, respectively. In the specific dynamical situation, the downward transport of the GMS component from the MLT is fast enough to reach the photolytic “safe” stratosphere at the end of the NH winter which then causes a significant NO_y enhancement in the upper stratosphere. It stays there and is diluted during the summer and is then transported further down in the beginning of the winter season. The enhancement related to the SPE is seen down to the tropopause and stays there over the whole year, with a relative enhancement of more than 20%. NO_x is lost through transport into the troposphere, where the main reservoir in the lowest stratosphere HNO_3 is washed out, or into the mesosphere, where it is photolytically destroyed.

In the SH (shown in Fig. 3), the NO_y enhancement is essentially from the SPE, as the MLT component is photolytically destroyed during the summer season. We also observe a higher relative enhancement here, as the excess NO_y in the upper stratosphere is transported downward in the beginning of winter compared to the upward draft in the NH summer. In addition, de- and renitrication by sedimentation of condensed HNO_3 can be observed, resulting in an effective removal of nitrate from the stratosphere into the troposphere. The relative maxima in the southern winter middle stratosphere are related to the higher supply of HNO_3 at locations in the reference run where this air mass already was depleted.

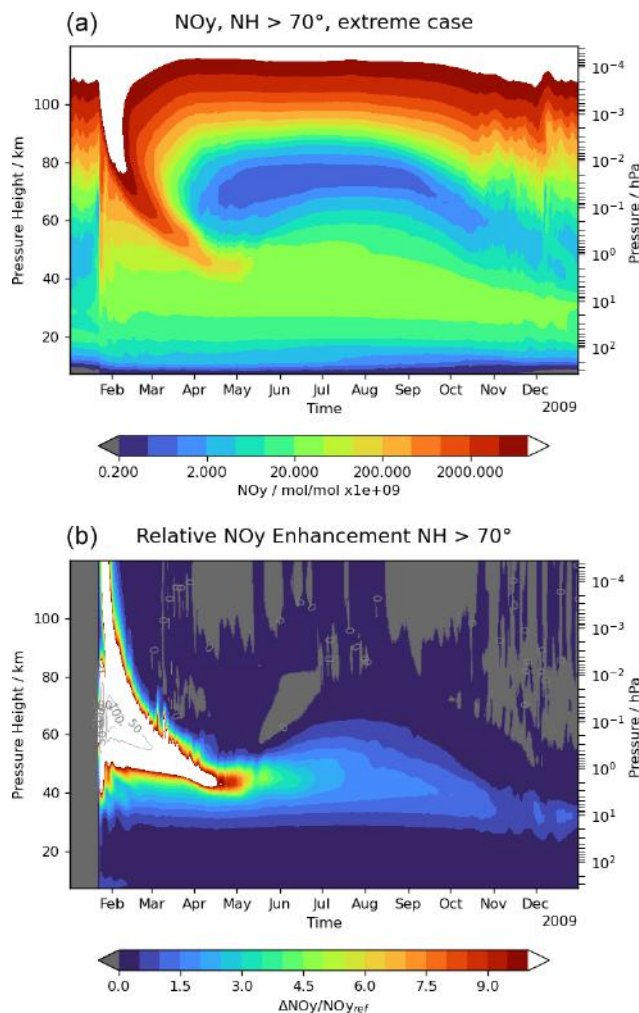


Figure 2. (a) NO_y mixing ratio in the course of 1 year with specified dynamics of the year 2009 (ERA-Interim) and applying the extreme scenario case as described in the text. For the polar cap, the intrusion connected to the MLT part dominates around the stratopause; below the stratopause NO_y stems from the SPE. (b) Relative enhancement of NO_y compared to the reference run.

The contribution of the GMS to the NO_y at high northern latitudes ($> 70^\circ \text{N}$) in the intrusion is shown in Fig. 4. It is limited in altitude to above about 40 km and is diluted to values below 0.1 in autumn.

The distribution of NO_y over latitudes at 37 km altitude is shown in Fig. 5. At this altitude the contribution of NO_y to ozone depletion maximizes (see, for example, Fig. 6.1 of Brasseur and Solomon, 2005). A closer inspection shows that at this altitude also NO_y from the NH is transported to the SH and significantly contributes to the NO_y after some months, as part of the diabatic circulation from the summer to the winter pole.

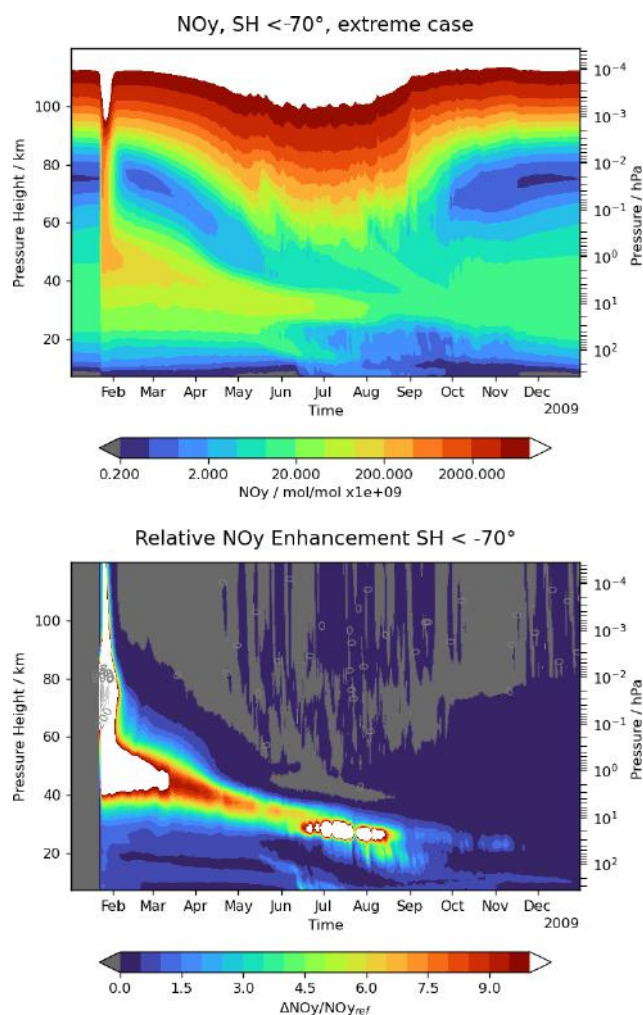


Figure 3. As in Fig. 2, but for 70° S.

4.2 Impact on the global NO_y budget

To assess the impact of the NO_y intrusions on the global NO_y budget, we first calculated the total column of NO_y with and without the event, as shown in Fig. 6. The two hemispheres show remarkable differences, with a higher total relative NO_y enhancement at high southern latitudes in the first 6 months but a fast decay after, in addition to a longer-lasting NO_y enhancement in middle and high latitudes in the NH (see also Fig. 7). From Fig. 3, which shows mixing ratios, one can deduce that the strong relative enhancement is affected by differences in the depletion of gaseous NO_y by denitrification in the reference and extreme run. This becomes clearer when showing the absolute change in the total columnar content. The change in the total amount NO_y (in units of mol) for different latitude bands by the event is shown in Fig. 7. In the SH, we clearly observe the fast decay shortly after the event connected to the photolytical destruction in the polar mesosphere and then some stabilisation in the winter followed by the denitrification in the cold polar stratosphere. The NH is

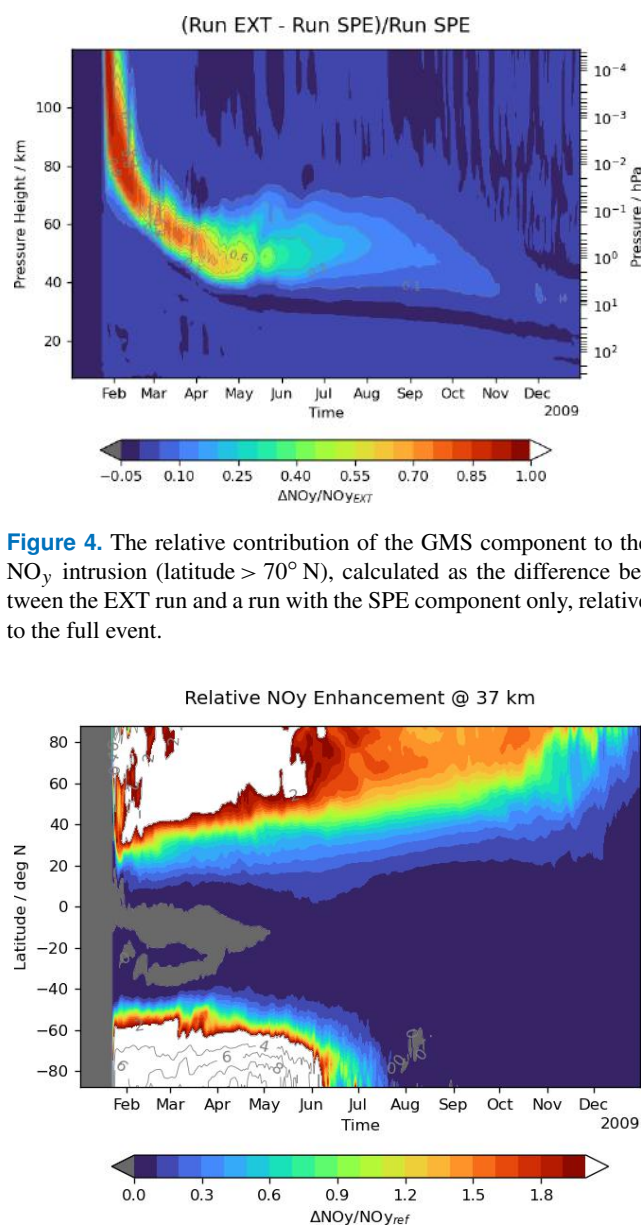


Figure 4. The relative contribution of the GMS component to the NO_y intrusion (latitude > 70° N), calculated as the difference between the EXT run and a run with the SPE component only, relative to the full event.

Figure 5. Relative enhancement of NO_y caused by the extreme event at 37 km pressure height.

characterized by a more continuous decay, stabilizing in the next winter season.

4.3 Impact on ozone

The NO_y enhancement causes additional ozone depletion via the catalytic reaction of NO_x with O₃. This is shown in Fig. 8 as the relative change in ozone. In the MLT, there is also a short-lived ozone enhancement which is connected to the dissociation of O₂ by atomic nitrogen. In addition, the ozone decrease in the stratosphere causes less absorption of solar UV and a small buildup of ozone below the depleted zone.

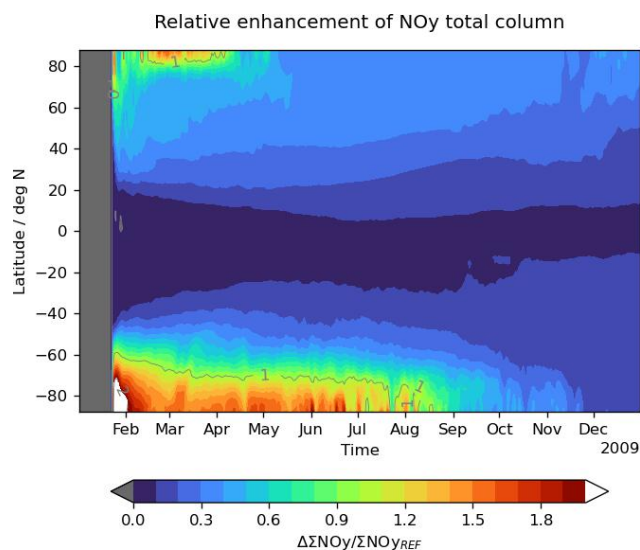


Figure 6. Enhancement of the total NO_y column (> 7 km pressure height) relative to the reference run.

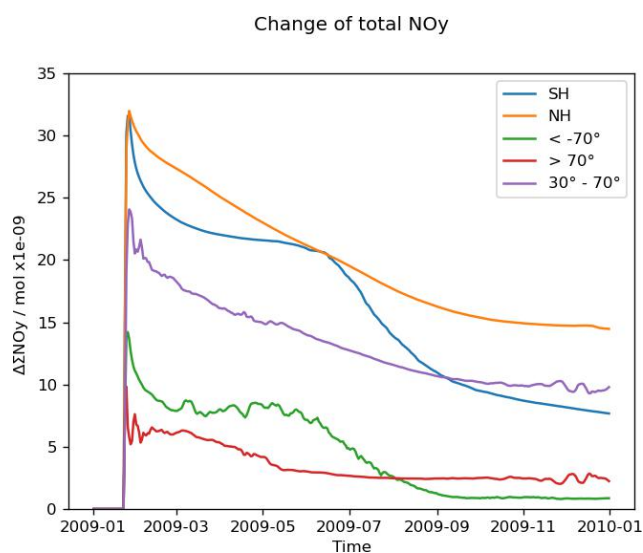


Figure 7. The NO_y signal from the extreme event (EXT-REF) for different latitude bands as the total amount of NO_y in moles.

Mirroring the different amount of NO_y transported into the stratosphere, the corresponding ozone loss is higher in the SH, whereas the ozone loss in the NH lasts for a longer time. This just reflects the loss of NO_y in the SH in the course of the polar winter.

The change in the ozone column relative to the reference run is shown in Fig. 9. It shows the previously described asymmetry in the two hemispheres, with exceptional high changes of total ozone loss in the Antarctic vortex. Total ozone recovers to nearly normal values at the end of the simulation year in the SH. On the other hand, in the NH the maximum loss in the mid-latitudes is less than 5%, but a deficit

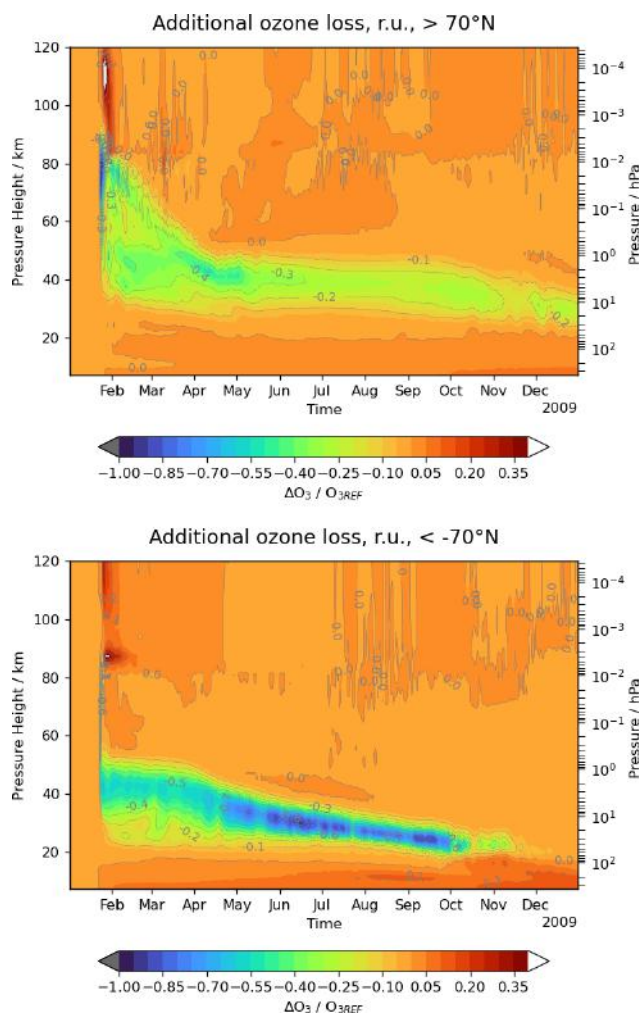


Figure 8. Ozone change in the extreme run relative to the reference run ((REF-EXT)/REF) related to the NO_y intrusion through the extreme event (EXT-REF) for the northern and southern polar caps (latitude $|\phi| \geq 70^\circ$).

of total ozone of a few percent stays permanently throughout the year. The contribution of the GMS to the loss in total ozone is negligible (not shown).

5 Discussion

5.1 Specific dynamical situation and the total NO_y input

Inspecting Fig. 1, we note that the ionization rate of the GMS shows its maximum in the thermosphere and at the auroral oval latitudes, whereas the extreme solar particle event (ESPE) is concentrated to the poles. This has some consequences for the respective impacts of the two components. As discussed in Sect. 2, the specific dynamical situation maximizes the impact of solar extreme events in the NH via long-lasting downward transport of NO_y in the winter polar middle atmosphere. Figure 4 shows that the contribution of the

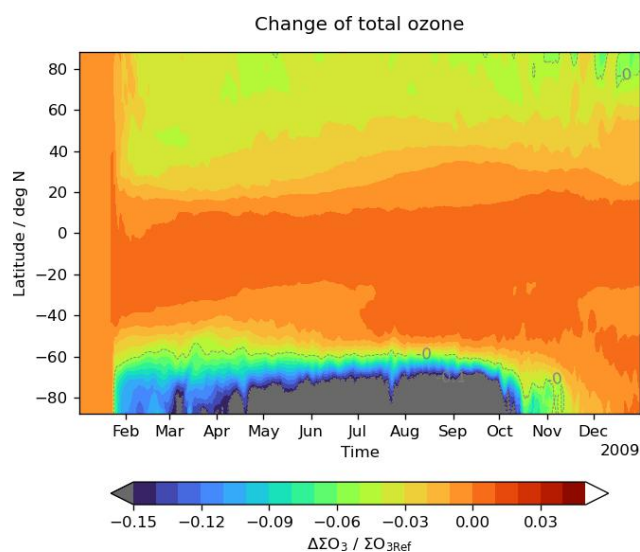


Figure 9. The change in total ozone caused by the solar extreme event relative to the reference run.

GMS to the additional NO_y is of minor importance compared to the impact of the ESPE. The higher altitude of the source region together with its lower latitude makes the photochemical destruction of NO_y at the end of the downward transport regime more probable. Accordingly, an additional experiment with a $10 \times$ GMS event, but without the ESPE component, was carried out. In this event, NO_y enhancement in the NH of about 2–3 Gmol in late summer is observed. We therefore conclude that an extreme GMS in the range discussed would not pose a strong and long-lasting disturbance in the NH middle atmosphere. Still, the mesospheric part of the ESPE benefits from this dynamical situation in the NH: NO_y that has not reached the mid-stratosphere would have been destroyed photochemically when transported by an unstable vortex or minor SSWs to the sunlit mid-latitudes.

For comparison, we also tested other dynamical situations, from early to late winter (only in the winter season is NO_y stable in the mesosphere): in winter 2009–2010 but the event put on DOY 300 in 2009, a rather early and strong mid-winter warming (December 1998), and the SSW in mid-January 2004, the date of which being quite similar to the setup studied. For the EXT scenario, we estimate the initial total NO_y input to about 65 Gmol which decreases to about 25 Gmol after 1 year (see Fig. 7). The undisturbed total NO_y amount in the atmosphere prior to the event above about 7 km altitude is simulated to 135 Gmol. The experiments for 2004 and 1998 yield the same initial amount of the intrusion but decay to about 10 and 25 Gmol after 1 year, respectively. The event set to DOY 300 in 2009 yielded, with 53 Gmol, slightly less than the initial NO_y and showed a rather fast decay to only a few gigamoles (Gmols) after 1 year.

Some insight might also be gained from comparison with Sukhodolov et al. (2017) (albeit their results are deduced

from an ensemble mean and for pre-industrial conditions): their Fig. 4 shows a decay of the NO_x enhancement to 100 % within about 140 d in the NH. In the EXT experiment after 140 d the enhancements still exceed 300 % and reach 100 % after about 270 d.

From these comparisons, the scenario selected seems to be an example of a strong and long-enduring impact of the event. Estimating the range of the impact of strong solar events under different dynamical situations requires additional simulations, including feedback effects on the dynamics. This is beyond the scope of this study and will be addressed in a follow-up paper.

The total NO_y mass produced in experiment EXT of about 65 Gmol, corresponding to about 40 % of the total atmospheric NO_y content, is substantially greater than what has been found in observation during the last decades. For example, Vitt and Jackman (1996) estimated a few percent increase of NO_y after strong SPEs in the northern polar stratosphere, Funke et al. (2014) estimated the total NO_y input from satellite observations for the period 2002 to 2012 to be a few gigamoles (Gmols) per year, and Reddmann et al. (2010) deduced the total NO_y input from a model simulation for the period 2003 to 2004 to be about 2 Gmol. In terms of NO_y input, the NO_y buildup by the solar event studied is about 10 to 30 times higher than what is estimated for solar particle events observed in the satellite era, which is on the order of the scaling factor we apply to the ionization rates of the SPE. In this sense, the solar extreme event is also an atmospheric extreme.

The budget of NO_y in the atmosphere in the longer term depends on the source strength of the event and the loss terms by photochemistry and exchange processes to the troposphere, where NO_y is removed by precipitation.

For the source strength, we use ionization rates estimated for the event in 774 CE. There are several steps to estimate the ionization rates for such events which all have factors of uncertainty. This is extensively discussed in Usoskin et al. (2011) and in Sukhodolov et al. (2017). One of the main uncertainties of the SPE part is the energy spectrum of the protons which has been scaled from an event in 1956 with a rather hard spectrum compared to other events in the satellite era Usoskin et al. (2020). A softer spectrum would shift the maximum NO_y production to higher altitudes, where the lifetime of NO_y is reduced. Therefore, the initial NO_y production by the ionization we estimate is probably on the upper limit of what can be expected.

The GMS contribution has been estimated from events observed in the past, but data at high flux levels are rare and have therefore substantial uncertainty. On the other hand, theoretical considerations (Vasyliunas, 2011) set an upper limit to the strength of a GMS in terms of Dst to about 2500 nT, which is probably outside of the estimations we used. In addition, the geographical distribution of the auroral oval depends on the strengths of the event. Here we used the distributions from the AISstorm model for observed events

and therefore miss the shift of the auroral oval to lower latitudes when the GMS strength is much higher. The scaling factors taken from Meredith et al. (2016) have been derived for a particle population that is partly trapped in the magnetosphere. Applying them to particle precipitation includes another source of uncertainty. But even a contribution greater by a factor of 2 from the GMS part would not bring this component above the SPE contribution.

For the efficiency of NO production, we use in the model the results of Porter et al. (1976). Some studies indicate a dependency of these factors on the ionization rate itself (Nieder et al., 2014), with a kind of saturation effect, especially in the thermosphere. This can be a major uncertainty for the production term and needs detailed ion chemistry studies.

Besides the source strength, the downward transport from the source region in the mesosphere to the stratosphere is essential. From the studies within the SPARC HEPPA initiative (Funke et al., 2011a, 2017), we know that many models do not correctly simulate this transport in the polar winter. The KASIMA model we use here in combination with ERA-Interim analyses yields downward transport from the MLT, which is generally in good agreement with observations, also for situations with an elevated stratosphere as we use here. Finally, the atmospheric lifetime of NO_y is also determined by the correct pattern and strength of the Brewer–Dobson circulation. It determines the upward transport into the region with photolytic loss and the exchange to the troposphere. The mean age of air simulations found with KASIMA is also in good agreement with observations (Stiller et al., 2008; Haenel et al., 2015).

In conclusion, the main uncertainties for the strength of the NO_y intrusion for a selected dynamical scenario as in our simulation are the energy spectrum of the particle flux and the initial ion chemistry production terms.

5.2 NO_y loss through the Antarctic winter

The loss of NO_y in the southern polar winter atmosphere deserves special attention. The KASIMA model includes heterogeneous processes in the polar winter stratosphere, causing denitrification and renitrication of the cold air masses in the Antarctic polar vortex through the nitric acid trihydrate (NAT) condensation of HNO₃ on water ice and the adsorption on sulfate aerosols, followed by sedimentation of larger particles into the lower layers as part of the annual cycle. In the KASIMA model, this process seems to be saturated: in the core of the Antarctic polar vortex, denitrification is complete. Additional NO_y brought into the lower stratosphere by the intrusion does not change this, but through the changed saturation pressure, greater volume will take part in denitrification. The change in total NO_y for the southern polar cap in Fig. 7 supports this finding: at the end of the winter, most of the additional NO_y of about 6 Gmol in the southern polar cap is lost to the troposphere.

Inspection of Fig. 3 shows that the majority of NO_y renitrication in the southern polar lower stratosphere takes place between the end of July and the middle of August. NO_y here is mainly in the form of HNO₃. Once transported to the troposphere, NO_y is deposited onto the surface by rain or snow within a few days. If we assume that the NO_y is deposited in the form of NO₃ on the southern polar cap, we deduce an average deposition value of $2.4 \times 10^{-6} \text{ g cm}^{-2}$. Using the approximation that the column-integrated particle produced NO_y is deposited, Melott et al. (2016) estimate a deposition of about $1.4 \times 10^{-7} \text{ g cm}^{-2}$ for the 1956 event, corresponding to $1 \times 10^{-5} \text{ g cm}^{-2}$ for the extreme event just by applying the scaling factor. As in our model, loss processes such as photolysis and transport to lower latitudes are included, and our lower value is in line with their estimation. With an average precipitation of 166 mm yr^{-1} for Antarctica and assuming that the nitrate is deposited within 1 month, we estimate a mass mixing ratio for nitrate of 190 ng g^{-1} in a monthly layer of (evenly distributed) deposited snow. A typical detection limit for nitrate events measured in ice cores with monthly-to-sub-annual resolution is on the order of 100 ng g^{-1} (Smart et al., 2014), so we would expect that a nitrate signal would be observable for such an event in Antarctica. Note that our model does not include tropospheric meteorology, so an analysis of regional precipitation patterns is outside the scope of this paper. Sukhodolov et al. (2017) estimate from their analysis that a nitrate signal would not be detectable in a Greenland ice core. This is essentially in agreement with our findings, which show a much smaller change in total NO_y in the northern polar cap in early spring. Here, the higher polar temperatures and less stable polar vortex do not allow for strong and fast denitrification, the primary cause of the high NO_y depletion in the southern polar lower stratosphere. In addition, we use analyzed winds and temperatures similar to realistic meteorology, whereas free-running models may fail in several ways to correctly describe the polar dynamical and chemical processes in the lower stratosphere. For the case of the 774/5 event, ice core data for Antarctica do not show nitrate enhancements (Sukhodolov et al., 2017; Jong et al., 2022). From Fig. 3 one may note that a strong denitrification signal is only expected when the NO_y tongue reaches the lower stratosphere when there are also low temperatures. A mismatch, for example an event occurring in Antarctic spring, would presumably result in a much broader and therefore undetectable nitrate peak.

5.3 Ozone loss caused by the event

Inspecting Fig. 8, one may note the nearly complete ozone loss immediately after the event in the polar mesosphere. This raises the questions of the extent that the ozone loss scales with the NO_y enhancement and if the loss is near saturation, at least during the first weeks after the event. For this reason, we performed an additional sensitivity experiment where we scaled the strength of the ionization rates of

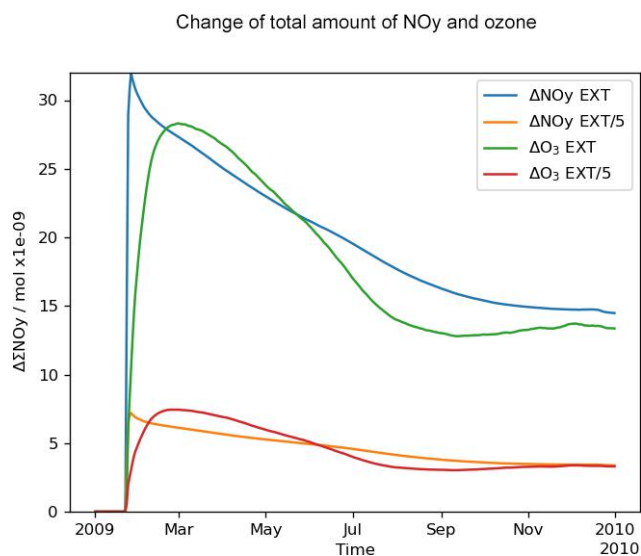


Figure 10. Total NO_y input by the EXT event and an event with a strength of 1/5 (in Gmols) into the NH. The corresponding decrease of the ozone (in Gmols), scaled by a factor of -0.03 , is also shown.

the EXT experiment by a factor of 0.2. Figure 10 gives the result of this experiment for the NH in terms of the total amount of the enhancement of NO_y and the corresponding ozone decrease. First, we can see that the total NO_y input into the atmosphere scales nearly perfectly with the strength of the forcing. Surprisingly, the ozone decrease also scales with the ionization rate by about the same factor, with a only slightly relative higher ozone decrease in the reduced case. In addition, the ozone loss is strongly correlated with the additional NO_y .

The reason for this behaviour lies in the fact that the ESPE reaches deep into the middle atmosphere, where the relative ozone reduction related to enhanced NO_y is small; again, see Fig. 2. For total ozone, which essentially stems from the lower stratosphere, this results in quasi-linear scaling of ΔNO_y and ΔO_3 .

Considering radiative feedback effects, the situation may be different: following Gray et al. (2010), an efficient top-down mechanism for radiative feedback involves temperature gradients induced by ozone changes in the mid-latitudes around the stratopause. For our setup, the induced decrease of the ozone amount above 40 km for latitudes $> 40^\circ\text{N}$ shows a more saturated scaling ($3\times$ instead of $5\times$) and vanishes after about 1 year, whereas the total ozone loss stays at a constant level in the second half of the year, stemming mainly from the lower stratosphere. In our setup, we would therefore expect only smaller feedback on the circulation compared to the setup of Sukhodolov et al. (2017).

5.4 Increase of UV index caused by the ozone loss

In order to estimate a health-related impact caused by the reduced total ozone column, we calculate the corresponding change in the UV index (UVI). The UVI is calculated from the spectrally weighted and integrated solar irradiance, where the weight function takes the erythemal action spectrum into account (CIE International Commission on Illumination, 2019). Here we use the simplified formula of Allaert et al. (2004), their Eq. (8) together with Eq. (6), to estimate the increase of UVI connected to the decrease of the total ozone column (ΣO_3) after a solar extreme event. In this formula, besides the dependence on the extraterrestrial flux and the solar zenith angle, for cloud-free situations, UVI depends on ΣO_3 only. For the column down to 7 km, the lower boundary of the model, we obtain an increase of the UVI for noon, as shown in Fig. 11. The absolute highest change of about 0.5 UVI occurs in the months of March and April at the latitude of 25°N , which includes the region on Earth with the highest population density. The relative change in the UVI reaches about 4%–5% in the northern mid-latitudes, where the summer months are also affected. In order to put this change in context, we remark that the ozone depletion by chlorofluorocarbons reached a maximum reduction of total ozone of about 3% in the northern mid-latitudes in 1995 (Weber et al., 2022), quite similar to the values we found for the extreme event (see Fig. 9) but lasting for more than a decade.

There is also a significant increase of UVI at high southern latitudes but not at latitudes which have high population density.

6 Conclusions

The purpose of this study was to analyze the direct effects of an extreme solar eruption on the chemical state of the middle atmosphere. We considered an extreme solar event with a strong SPE that was also connected with a strong GMS. The setup of the numerical experiment was chosen to yield a maximum chemical impact in the NH. This was achieved by setting the event to NH winter in combination with an elevated stratosphere SSW event. The main findings are that the GMS component from energetic electrons even in this dynamical situation leads only to a minor ozone reduction. The ozone reduction in our extreme scenario stems essentially from the ESPE and survives through the following boreal summer. The additional NO_y amount in the middle atmosphere scales with the strength of the event, as does the ozone loss. The studied event yields a significant but limited enhancement of the UVI in populated latitudes in the first months following the event of a few percent. The ESPE causes a strong impact in terms of additional NO_y and ozone loss in southern high latitudes, despite the event having been put to NH winter to minimize photochemical loss of NO_y . This is related to the deep penetration of energetic particles at altitudes where the

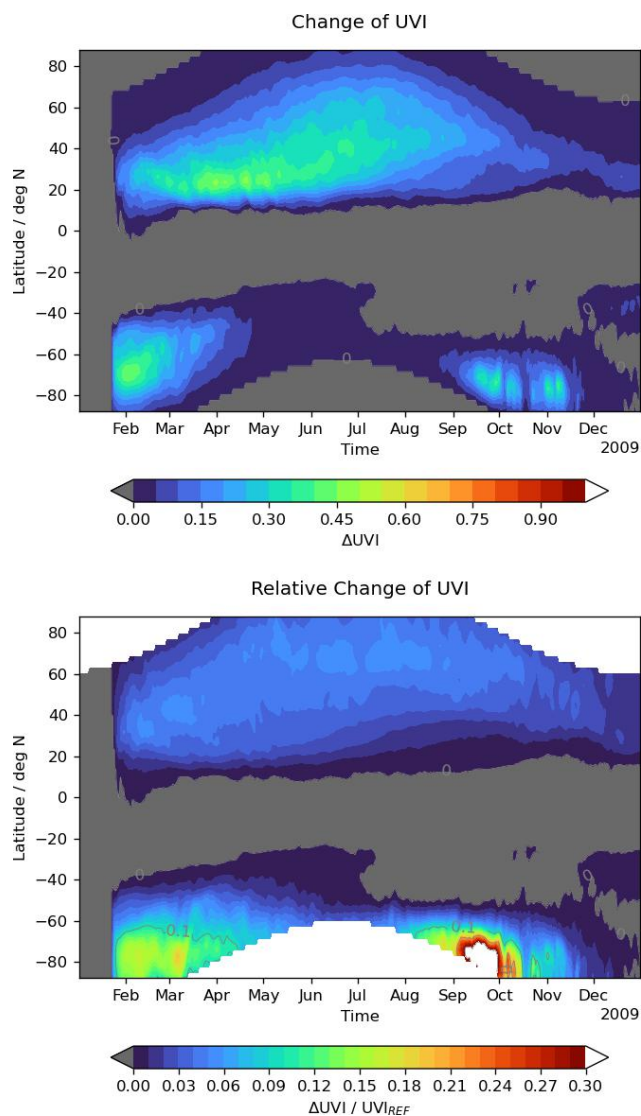


Figure 11. Change and relative change in the UVI caused by the extreme event for noon.

loss of NO_y is negligible and also to the early onset of downward transport with the beginning of Antarctic winter in the middle atmosphere. Besides the related ozone loss, the denitrification of NO_y in the lower Antarctic stratosphere and the following washout in the troposphere can yield potentially measurable NO_3 signals in ice cores. This, however, has not been observed to date, perhaps because the signal strength strongly depends on the timing of the event. Therefore, nitrate signals in ice cores are not reliable indicators of solar activity.

Data availability. Model data used for this analysis are accessible through <https://doi.org/10.35097/1104> (Reddmann, 2023).

Author contributions. TR prepared the experiments and wrote the paper; IU, JMW, and OY provided the ionization rates; and MS critically discussed and improved the manuscript.

Competing interests. The contact author has declared that none of the authors has any competing interests.

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